


UMR Gulliver


Effets Collectifs & Matière Molle

Jamming, (8 years of) a granular media experiment


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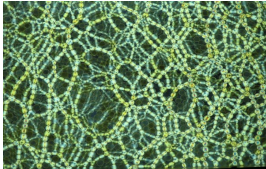

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Cargèse 28/08/2014

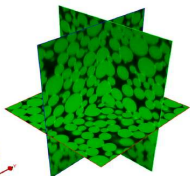
Experimental realizations



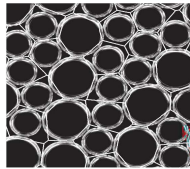
Green peas, Hales, 1727





Grains, Behringer



Emulsion, Jorjadze et al., 2011

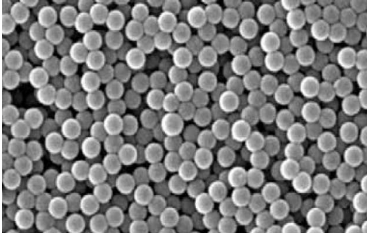


Foam, Katgert et van Hecke, 2010

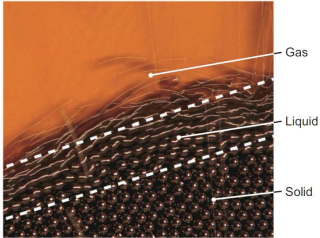


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What about *these* situations?





Colloidal suspensions
=> thermal agitation



Dense granular flows
=> mechanical excitation

Control of Dynamics by Jamming scalings?

Effect of Dynamics on jammed systems?



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Initial goal!

Consider a (very) gently vibrated system of grains and study:

The glass and the jamming transitions

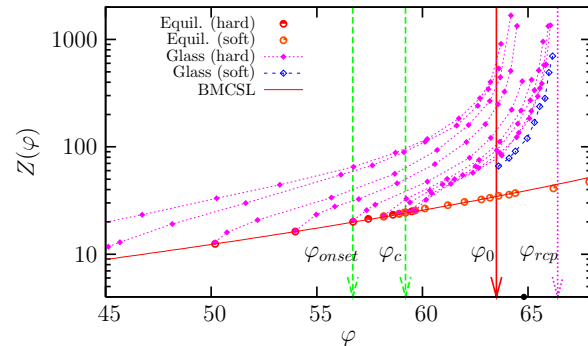
- Spontaneous fluctuations : vibrational dynamics vs. relaxation
- Shear Modulus : linear vs. non-linear regime
- Yield stress : thermal vs. mechanical origin
- Flows : rheology



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What is the plan?

- Reach jamming, i.e. enter deep into the glass phase



[Chaudhuri *et al.*, PRL '11]

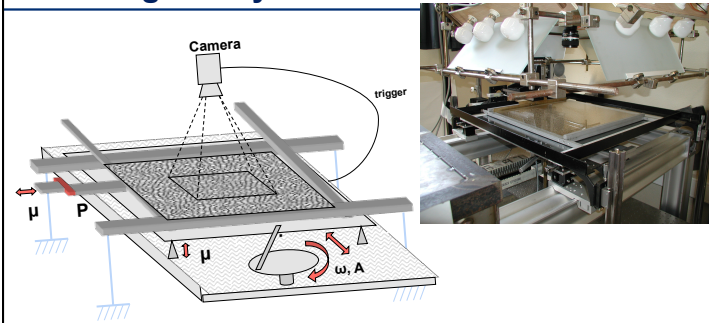


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Jamming in a system of vibrated brass discs



- Horizontal vibration ($\omega=10$ Hz, $a=1$ cm)
- Bi-disperse : $d_s = 4$ mm $d_l = 5$ mm
- 8000 brass discs in the system (1500 tracked)
- Vibration-triggered camera
- Tunable volume
- Pressure measured on the side



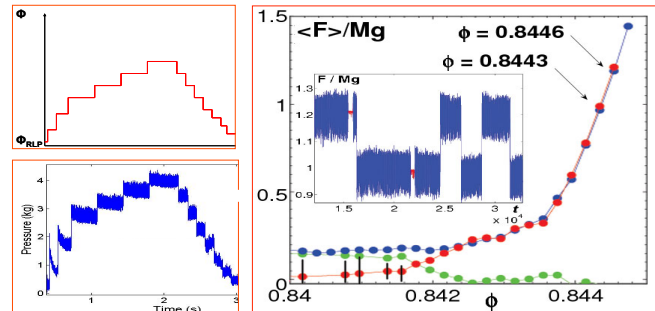
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Experimental protocol

- Increase packing fraction stepwise:
 - Allow for the slow relaxation of pressure
- Then decrease packing fraction and record dynamics



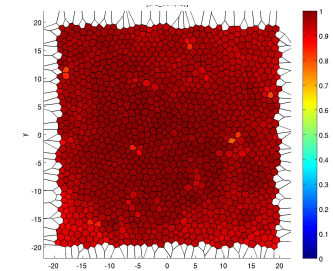
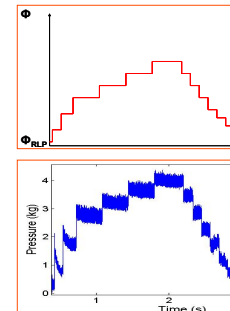
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Experimental protocol

- Increase packing fraction stepwise:
 - Allow for the slow relaxation of pressure
- Then decrease packing fraction and record dynamics



A completely frozen structure
⇒ A granular glass



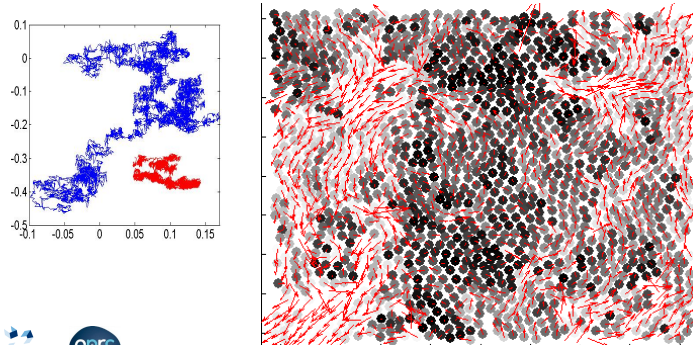
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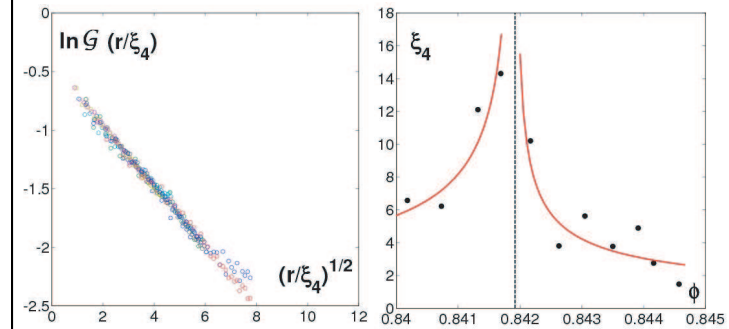
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Dynamics: Heterogeneous tiny displacements

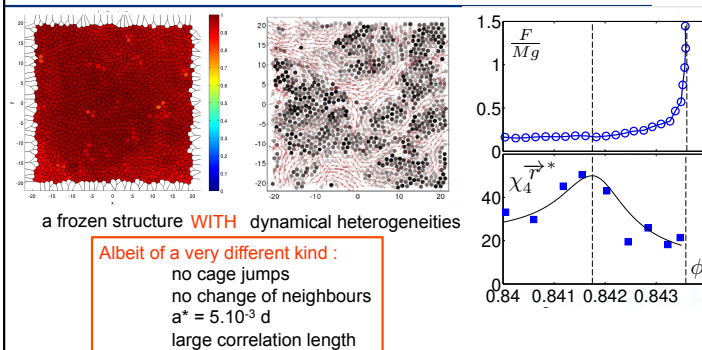
- Particles trajectories : $\vec{r}_i(t)$
- Displacement : $\Delta\vec{r}_i(t, \tau) \equiv \vec{r}_i(t + \tau) - \vec{r}_i(t)$



Dynamical heterogeneities

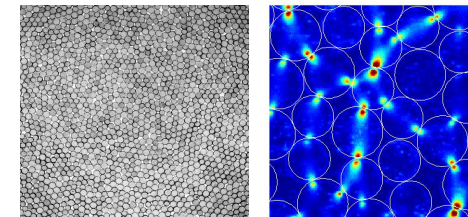
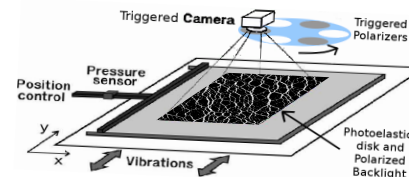


Altogether...



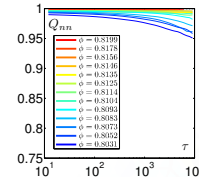
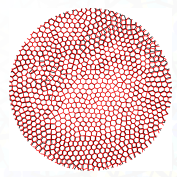
- What is the mechanism responsible for such heterogeneities?
- Why is there a maximum and not just a divergence, with the pressure?

Redo the experiment with soft photoelastic discs => access to contacts

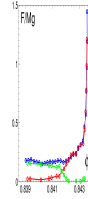
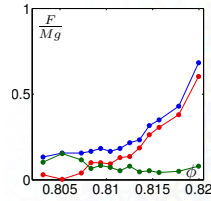


Same protocole: again a granular glass

- A frozen structure

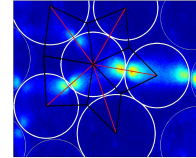


- But this time a glass of **soft** discs

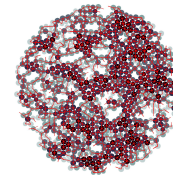
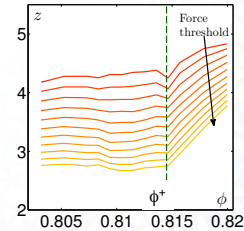


Signature of jamming within contacts

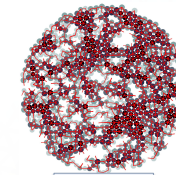
Interparticle force measurement



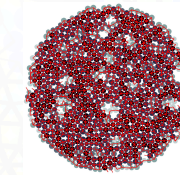
thresholding
gap < ε
→
thresholding
force > f₀



$\phi < \phi^*$



$\phi^* = 0.814$

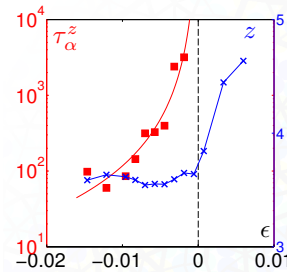
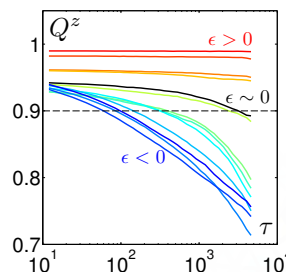


$\phi > \phi^*$

Dynamics of the contact network...

$$Q^z(t, \tau) = \frac{1}{N} \sum_i Q_i^z(t, \tau) \quad \text{where} \quad Q_i^z(t, \tau) = \begin{cases} 1 & \text{if } |z_i(t + \tau) - z_i(t)| \leq 1 \\ 0 & \text{if } |z_i(t + \tau) - z_i(t)| > 1 \end{cases}$$

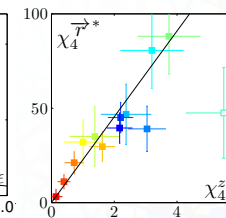
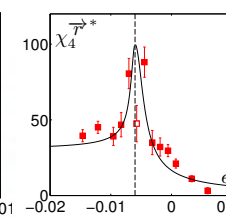
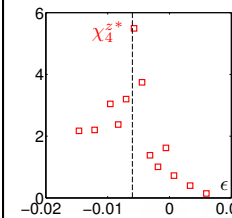
$$Q_z(\tau) = \langle Q^z(t, \tau) \rangle_t$$



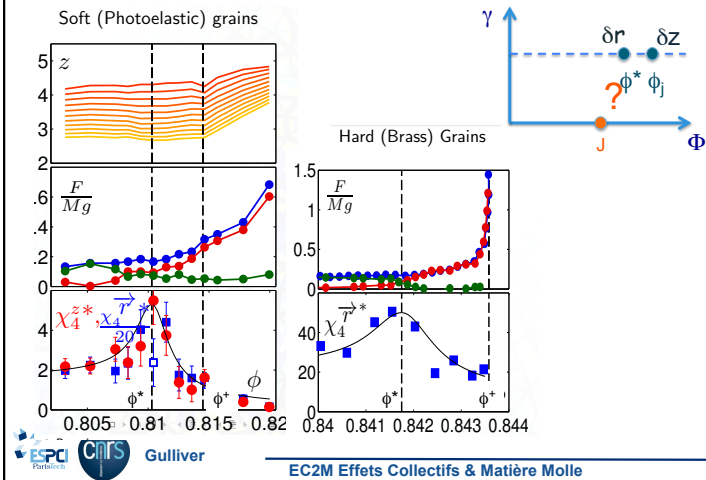
Dynamical arrest of the contact dynamics...

...is heterogeneous and governs the grains motion

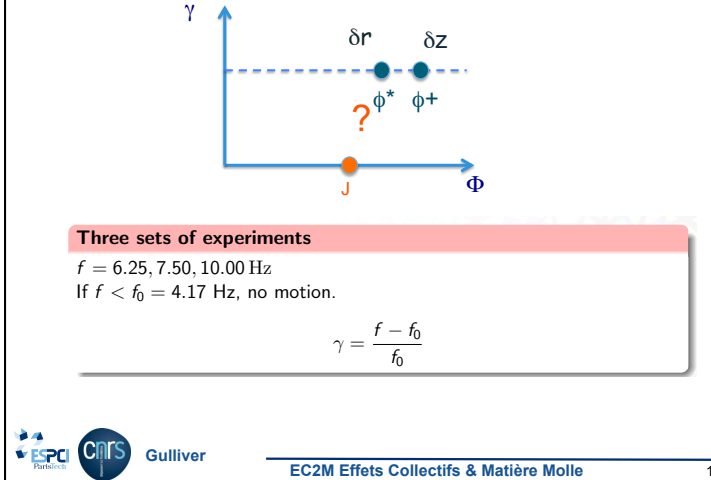
$$\chi_4^{z*}(\tau) = N \text{Var} \left(\langle Q_i^{z*} \rangle_i \right)_t \quad \left\{ \begin{array}{l} Q_i^z(t, \tau) = \begin{cases} 1 & \text{if } |z_i(t + \tau) - z_i(t)| \leq 1 \\ 0 & \text{if } |z_i(t + \tau) - z_i(t)| > 1 \end{cases} \\ Q_i^z(t, \tau) = \exp \left(-\frac{\Delta r_i^2}{2 \langle \Delta r_i^2 \rangle} \right) \end{array} \right.$$



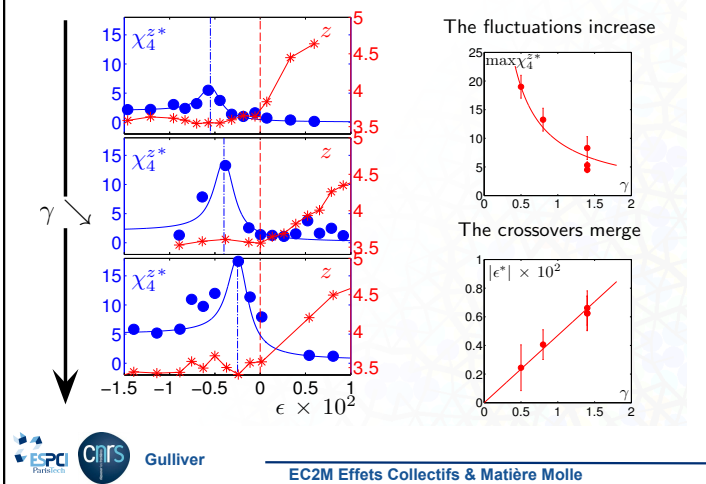
Summary: two distinct signatures



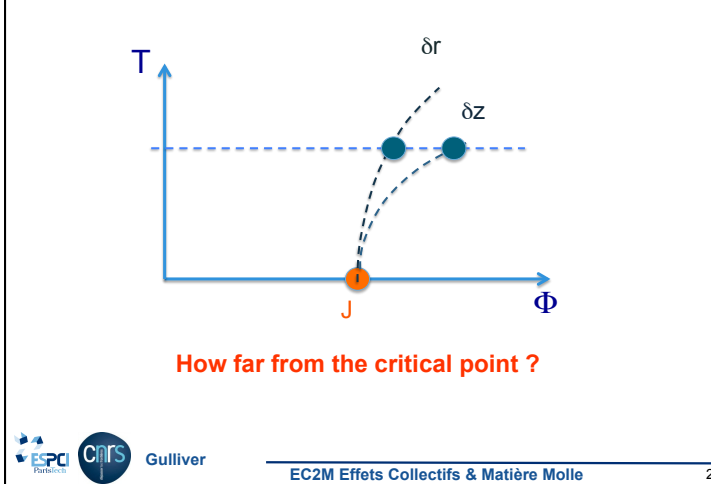
Reducing the vibration



Decreasing the vibration



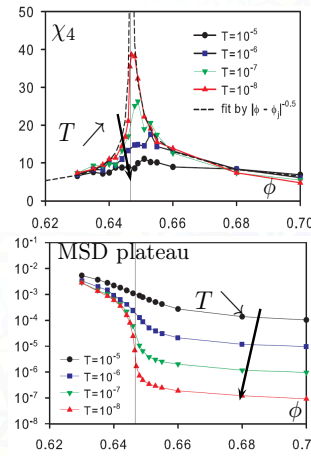
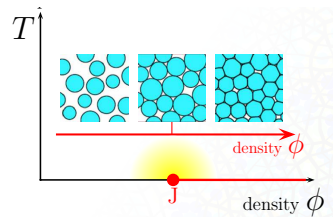
Hence two crossover lines



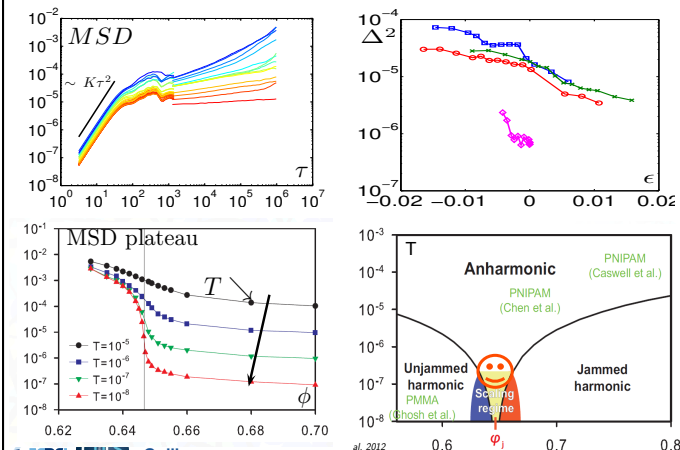
Comparison with thermal soft spheres...

Simulation of **thermal soft-spheres**

Ikeda et al, 2012



Comparison with soft spheres

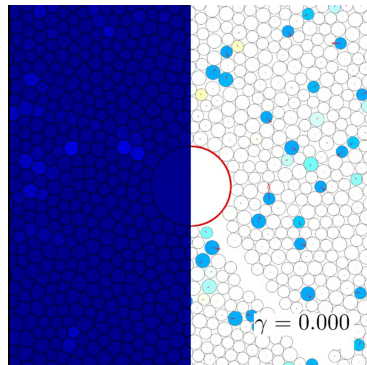


Probing elasticity

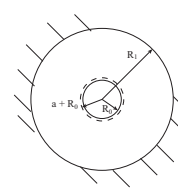
- Prepare the system at large packing fraction under vibration
- Inflate an intruder in the center (the vibration is stopped)
- Decrease the packing fraction while vibrating
- iterate

$$R_0 \rightarrow R_0 + a$$

$$\gamma = a/R_0$$



Probing elasticity : the linear elastic framework



$$\text{div}(\underline{\underline{\sigma}}) = 0$$

$$\underline{\underline{\sigma}} = \frac{1}{2} \text{Tr}(\underline{\underline{\sigma}}) \underline{\underline{1}} + \underline{\underline{\tau}}$$

$$\underline{\underline{\sigma}} = K \text{Tr}(\underline{\underline{\epsilon}}) \underline{\underline{1}} + 2G \underline{\underline{\tau}}$$

$$\underline{\underline{\epsilon}} = \frac{1}{2} [\underline{\underline{\nabla U}} + {}^t \underline{\underline{\nabla U}}] = \frac{1}{2} \text{Tr}(\underline{\underline{\epsilon}}) \underline{\underline{1}} + \underline{\underline{\gamma}}$$

$$U(R_0) = a$$

$$U(R_1) = 0$$

$$\delta = \text{Tr}(\underline{\underline{\epsilon}}) = -2 \frac{a}{R_0} A$$

$$\gamma = J_2(\underline{\underline{\gamma}}) = \sqrt{\frac{1}{2} \underline{\underline{\gamma}} \circ \underline{\underline{\gamma}}} = \frac{a}{R_0} B \left(\frac{R_0}{r} \right)^2$$

$$P = \text{Tr}(\underline{\underline{\sigma}}) = K \text{Tr}(\underline{\underline{\epsilon}})$$

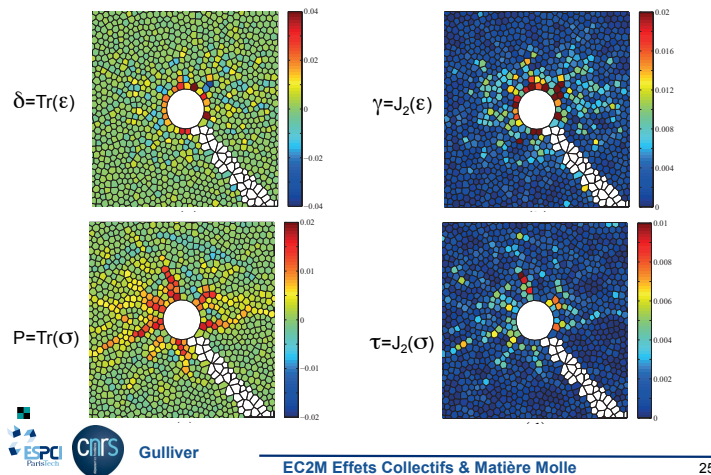
$$\tau = J_2(\underline{\underline{\tau}}) = \sqrt{\frac{1}{2} \underline{\underline{\tau}} \circ \underline{\underline{\tau}}} = 2G J_2(\underline{\underline{\gamma}})$$

$$A = \frac{R_0^2}{R_1^2 - R_0^2}; B = \frac{R_1^2}{R_1^2 - R_0^2}$$

Nota Bene

- In the limit of large R_1 , $A \rightarrow 0$, $B \rightarrow 1$: this is a shear test!
- G and K are simply obtained by the ratio of the stress and strain tensor invariants

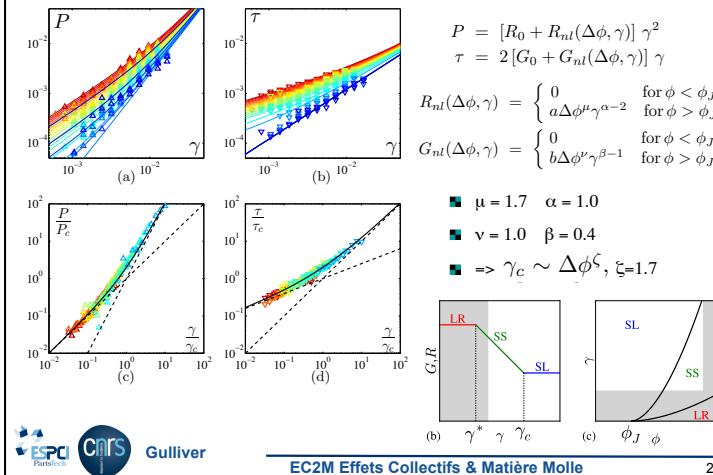
For each packing fraction and each a/R_0



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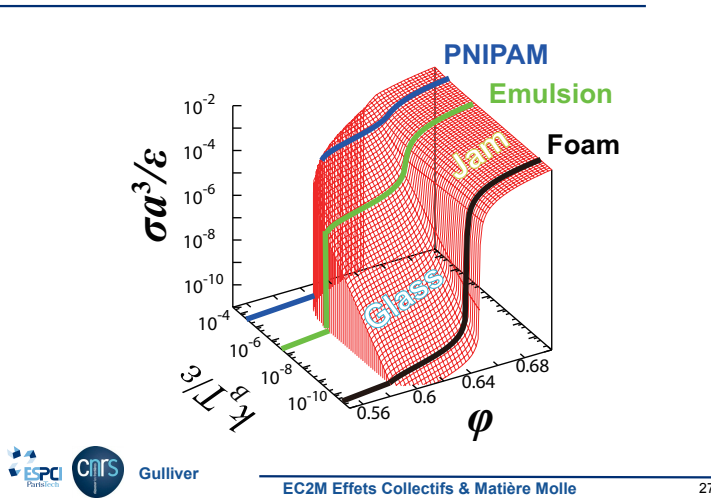
Parametric plot of stress vs strain



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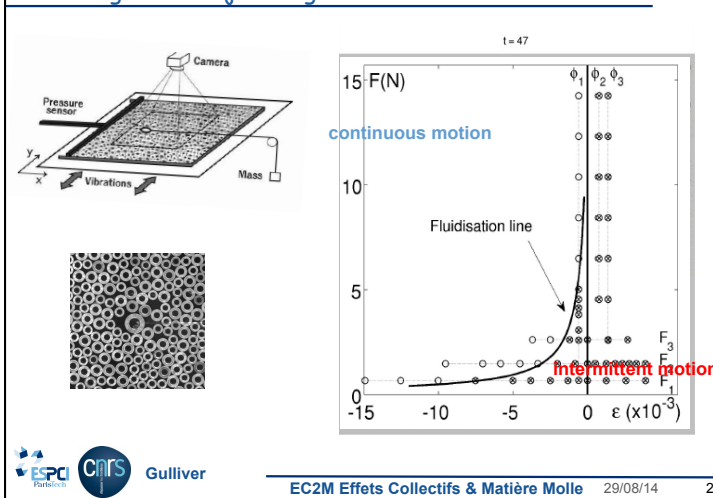
Yielding close to jamming ...



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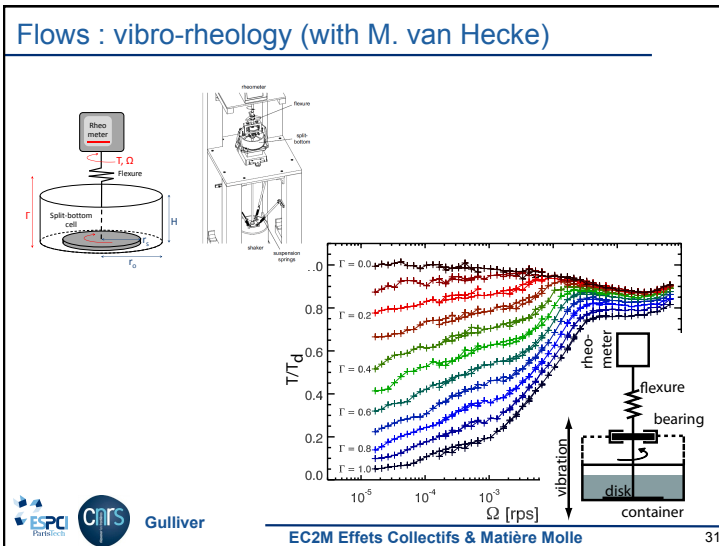
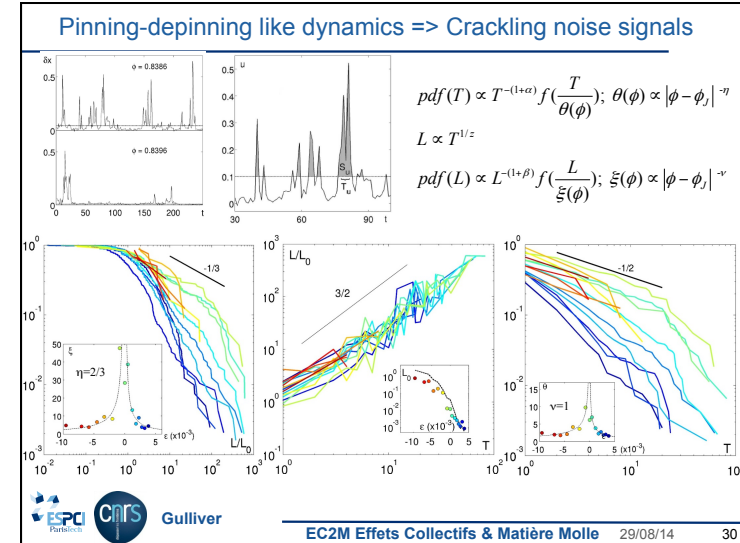
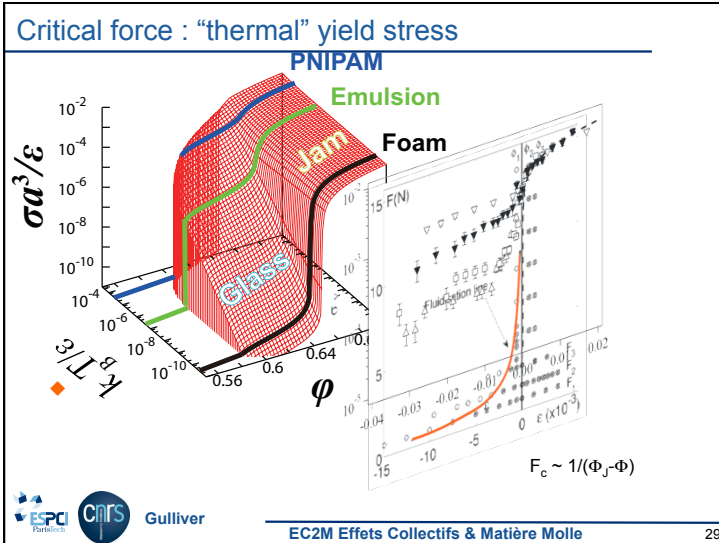
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Yielding close to jamming : the motion of an intruder ...



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Conclusion

- Vibrated granular media are suitable tools for probing the vicinity of jamming, (in particular low enough T_{eff})
- Two distinct crossovers (one dynamical, one structural) converge toward J-point in the limit of low vibration
- Inflating an intruder in soft photo-elastic discs => Non linear rheology
- Pulling an intruder in vibrated hard discs => the yield stress of “thermal origin” and reveals complex pinning – depinning like dynamics
- Vibro-rheology : flow curves close to jamming
- **Thank you!**

Further readings :

- Europhysics Letters, 83, 46003, (2008).
- Soft Matter, 6 (13), 3059–3064, (2010).
- Phys Rev Lett 103 12800 (2009).
- Europhysics Letters, 100, 44005 (2012).
- Soft Matter 10 (10), 1519-1536 (2014).
- Phys Rev Lett (2014) to appear